

Axial Crushing of Concentric Expanded Metal Tubes Under Impact Loading

Abstract

This paper investigates the axial crushing response of concentric expanded metal tubes under impact loading. The study is conducted by means of nonlinear finite element simulations, considering the speed of deformation and the size of the expanded metal cells. In the analysis, two tubes are placed concentrically and then axially crushed with an impact mass at various speed levels. Crushing load-displacement responses are analyzed to investigate the influence of the size of the mesh and the impact speed. The results show that specimens with the same mesh sizes increases their crushing length with increasing speed. Finally the main outcome of this investigation is to add to the state of knowledge information concerning the use of expanded metal meshes in impact attenuation devices.

Keywords

Impact loading, axial crushing, expanded metal, collapse modes.

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1 INTRODUCTION

Expanded metal meshes have been widely used in the industry for decorative purposes, protective and staircase systems. In spite of the growing number of patents developed every year (Smith et al. 2009a), research projects concerning their structural behavior are still limited. As seen in previous works (Graciano et al. 2009, Smith et al. 2014), tubes fabricated with expanded metal meshes are feasible alternatives for energy absorption applications due to their collapsibility.

Commonly, lightweight structural members subjected to compressive loads exhibit good collapsibility, nevertheless this feature is compromised by geometric and material imperfection sensitivities that reduces their carrying capacity. To guarantee a stable and progressive collapse the members need enough material availability for plastic deformation. It can be achieved by adding more tubes, while

keeping the thickness small. Therefore, placing or nesting tubes concentric or eccentrically is also a feasible alternative to enhance the performance of energy absorbers (Haghi et al. 2013). Lately, Nia et al. (2015, 2016a, 2016b) investigated the energy absorption characteristics of nested multi-tubular structures. Inserting foam between the tubes modify the deformation patterns and smooth the unstable responses (Zhang et al. 2010, 2012; Niknejad and Orojloo 2016).

Nested tubes may be placed transversally, in inventive configurations, for crashworthy applications (Morris et al. 2006, 2007; Baroutaji et al. 2016). By placing the tubes transversally, the energy absorption mechanism changes from axial crushing to a bending mechanism. Several assets are achieved using these type of tubes, for instance the possibility of nesting an array of tubes of various diameters reducing the occupied space, and furthermore the tubes may undergo additional deformation to obtain oblong configurations which allow larger crushing displacements than for circular ones (Baroutaji et al. 2014).

Martínez et al. (2013) investigated numerically the energy absorption capacity of concentric expanded metal tubes using the same cell orientation. Afterward, Smith et al. (2014) conducted an experimental investigation on the quasi-static response of concentric expanded metal tubes with various cell orientations, the results showed that the deformation pattern of the tubes changes significantly due to interaction between collapse modes for the inner and outer tubes.

Recently, Nouri et al. (2015) performed an experimental and numerical study on the response of expanded metal absorbers subjected to low speed impact. Subsequently, Hatami and Nouri (2015) also conducted a study on the energy absorption characteristics of multi-walled expanded metal energy absorbers considering a single cell orientation. Abdelaal and Tarlochan (2015) investigate the response of the tubes for crashworthiness applications. Graciano et al. (2016) studied numerically the effects of the elastoplastic behavior of the material and impact speed on the response of expanded metal tubes. Furthermore, combining design of experiments techniques and finite element simulations, Borges et al. (2016) studied also the impact response of these tubes. Summing up, the outcomes of these investigations show that the size and orientation of the expanded metal cells, and the impact velocity are the most influential parameters affecting the energy absorption response of expanded metal tubes.

As seen previously in the literature review, the behavior of expanded metal tubes in impact scenarios still requires further investigation. Therefore, this paper aims at investigating numerically the axial crushing response of concentric expanded metal tubes under impact conditions. Then, a parametric analysis considering the size and orientation of cells as well as the impact velocity is performed using nonlinear finite element modeling. The numerical procedure for the impact scenario is validated with previous studies taken from the literature.

2 FINITE ELEMENT MODELLING

2.1 Numerical Modeling for Expanded Metal Tubes

Expanded metal sheets are manufactured in a single process upon in-line expansion of partially slit metal sheets producing diamond like cells as shown in Figure 1 (EMMA 2015). These cells are composed by strands and connecting nodes. Figure 1 shows a schematic view of an expanded metal cell,

in which w is the strand width; t is the strand thickness, L_2 and L_1 are the minor and major cell axes, respectively.

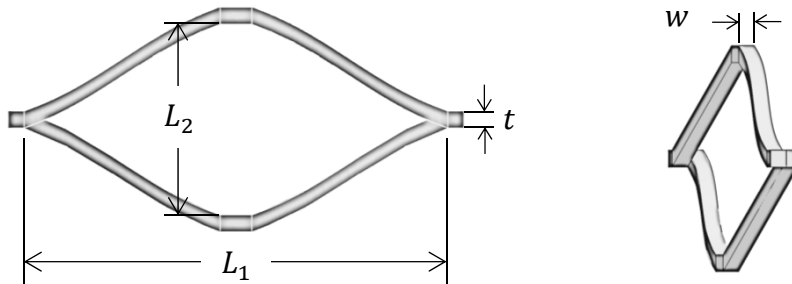


Figure 1: Expanded metal cell geometry.

To predict the axial crushing response of the concentric expanded metal tubes subjected to impact loading, the explicit nonlinear finite element code ANSYS AUTODYN (2010) is used herein. Individual cells in the tubes are modeled using 4-nodes tetrahedral solid elements (SOLID 164) suitable for modeling plasticity and large strain analyses. Figure 2 shows a typical mesh for an expanded metal cell. The contact between the rigid plates and the tube, and the self-contact between the strands of expanded metal as the tube collapses are taken into account using a penalty method based algorithm (ANSYS 2009). To simplify the numerical model all surface contacts in the model were treated as frictionless.

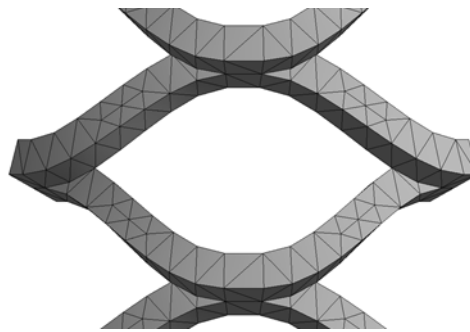


Figure 2: Mesh for an expanded metal cell.

The material used to model the expanded metal tubes is an ASTM A-569 steel, hence the following material properties are considered: Young's modulus $E= 205\text{GPa}$ and Poisson ratio $\nu= 0.3$; quasi-static yield stress $\sigma_y= 250\text{MPa}$; hardening coefficient $k= 495.35\text{MPa}$; and strain-hardening coefficient $n=0.21$. These values were obtained experimentally in (Sanchez 2005). An elastoplastic model was used to reproduce the steel properties of the expanded metal cells. In the elastic range, the material is assumed homogeneous, linear and isotropic, and in the plastic range, isotropic hardening is taken

into account. Therefore, the hardening curve of the material can be described employing the Hollomon correlation:

$$\sigma_p = k \cdot \varepsilon_p^n \quad (1)$$

In Eq. (1), the plastic stress (σ_p) is expressed as a function of the plastic deformation (ε_p), the hardening coefficient k and the hardening exponent n . Furthermore, the effect of the strain rate of the material is taken into account in the numerical model by implementing the Cowper - Symonds correlation (Cowper and Symonds 1957) shown in equation

$$\varepsilon_p = D \cdot \left(\frac{\sigma_y^D}{\sigma_y} - 1 \right)^q \quad \text{where} \quad \sigma_y^D \geq \sigma_y \quad (2)$$

From Eq. (2), σ_y^D is the dynamic yield stress at a uniaxial plastic deformation rate $\dot{\varepsilon}_p$, σ_y is the static yield stress and D and q are material constants. For low carbon steels such as the one investigated herein, the corresponding values for $D=6844s^{-1}$ and $q=3.91$ can be used (Abramowicz and Jones 1984). Hence, the final plasticity model, considering Eqs. (1) and (2), is provided in Eq. (3), which expresses the magnitude of the stress or yield surface for the isotropic hardening of the material

$$\sigma_y^D = (\sigma_y + k \cdot \varepsilon_p^n) \cdot \left[1 + \left(\frac{\dot{\varepsilon}_p}{D} \right)^{\frac{1}{q}} \right] \quad (3)$$

Figure 3 shows a schematic view of the boundary conditions considered in the numerical model. Accordingly, the nodes at the bottom of the tube are fully clamped, while the nodes on the impacted end move only in the vertical direction allowing axial crushing. The impact mass m has an initial speed V towards the impacted end of the tube, which is initially stationary. The impact mass and the base of the tube are modeled as rigid square plates (180x180mm), using 4-nodes hexahedral solid elements, with an infinite stiffness, in order to account mass inertial effects only. In the numerical analyses, all the concentric tubes were constrained to an outer tube with a mesh oriented at $\alpha= 0^\circ$ (Figure 4a), and an inner tube with a mesh oriented at $\alpha= 90^\circ$ (Figure 4b). Tube orientation is referred to the direction of the major axis of the cells.

The numerical procedure used herein has been validated using the results from the investigation performed by Nagel and Thambiratnam (2004) on the dynamic response of tapered thin-walled rectangular tubes. For details, the reader are encouraged to read about the impact response of single expanded metal tubes in Borges et al. (2015, 2016). In this work, a further convergence analysis was performed for tubes with meshes H23A (Table 1) subjected to an impact mass travelling at a speed $V=10$ m/s. The final mesh has 10230 elements, 3585 with an average size of 5.4 mm for the outer tube, and 6645 with an average size of 7 mm for the inner one.

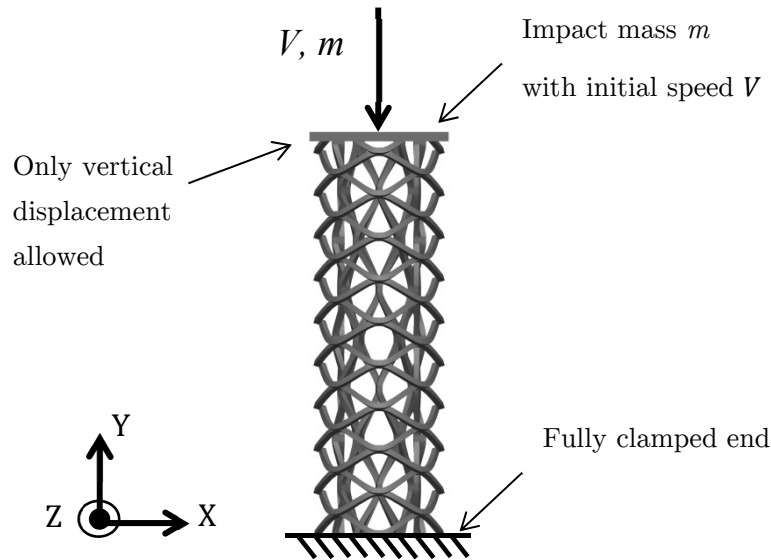


Figure 3: Schematic view and boundary conditions used in the numerical model of the concentric tubes.

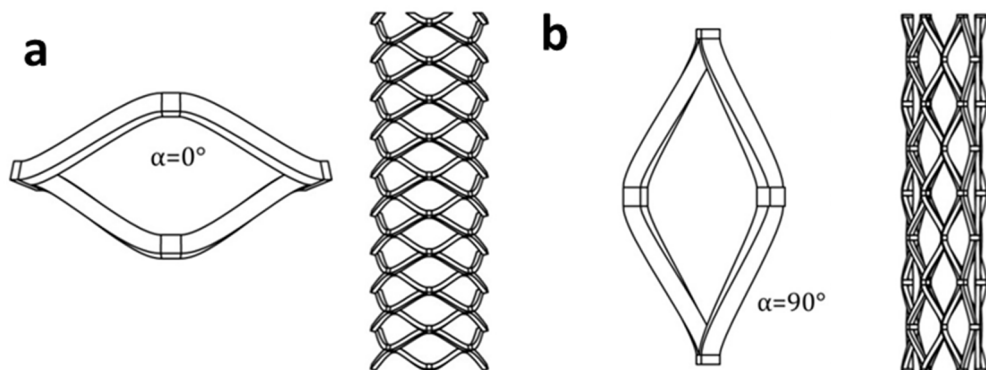


Figure 4: Cell orientations: (a) Outer tube $\alpha=0^\circ$; (b) Inner tube $\alpha=90^\circ$.

2.2 Parametric Analysis

Table 1 shows the dimensions of the three expanded metal cells analyzed herein. The finite element models were built according to the cell orientation, for $\alpha=0^\circ$ the models have 4 cells in the radial direction and for $\alpha=90^\circ$ the models have 6 cells. The length of the tubes is approximately three times their diameter.

Using the dimensions of the cells presented in Table 1, six numerical models (one for each single tube) were built to obtain three concentric tubes, *i.e.* 3 for the outer and 3 for the inner tubes. Each concentric tube was created using the same mesh type for outer and inner tubes. Table 2 shows the corresponding dimensions of the tubes, in which L is total length of the tubes, ϕ_e and N_e are the diameter and number of cells of the outer tube, respectively. Correspondingly, ϕ_i and N_i are the diameter and number of cells of the inner one.

Mesh	L_1 (mm)	L_2 (mm)	t (mm)	w (mm)
H23	64.5	30.6	4.1	4.5
H27A	81.6	36.8	6	6
H23A	58.6	29.2	7.5	6

Table 1: Dimensions of the expanded metal cells.

Tube	ϕ_e (mm)	ϕ_i (mm)	N_e (mm)	N_i (mm)	L (mm)
H23	115	73	4	6	409
H27A	109	80	4	7	407
H23A	134	84	4	6	428

Table 2: Dimensions of the concentric tubes.

It is important to notice that the tubes with $\alpha = 90^\circ$ were placed inside the tubes with $\alpha = 0^\circ$. In previous studies (Smith et al. 2014, Hatami and Nouri 2015), it was found that tubes with $\alpha = 90^\circ$ fail due to barreling of the cross section, i.e. outward buckling of the tube walls. Therefore, outer tubes with $\alpha = 0^\circ$ restrict the buckling deformation of the inner tubes, enhancing both the load-carrying capacity and the energy absorbed of the concentric tubes. All tubes underwent impact loading using a mass $m=100$ kg, and the impact speed was varied from $V=10$ m/s to $V=30$ m/s. Figure 5 show a plan view of the arrangement for the concentric tubes.

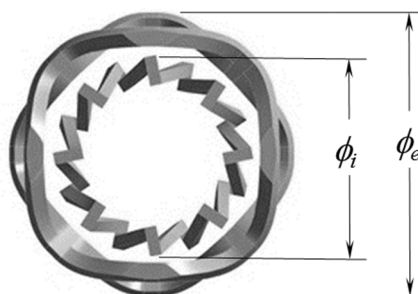


Figure 5: Plan view for the concentric tubes.

3 RESULTS AND DISCUSSION

Figures 6 to 8 shows the load-displacements responses obtained numerically for the concentric tubes described in Table 2 as a function of the impact speed V . From these curves, the following parameters were calculated: mean load (P_m) and the absorbed energy (E_a), and the crushing length (l_c). All these parameters are calculated up to a crushing length l_c that varies for each geometry.

The mean crushing force P_m , based on the area under the curve over the crushing distance l_c , is calculated by

$$P_m = \frac{\int_0^{l_c} P dl}{l_c} \quad (4)$$

Then, the absorbed energy is calculated integrating the load-displacement curves as

$$E_a = \int_0^{l_c} P dl \quad (5)$$

where P is the measured crushing force.

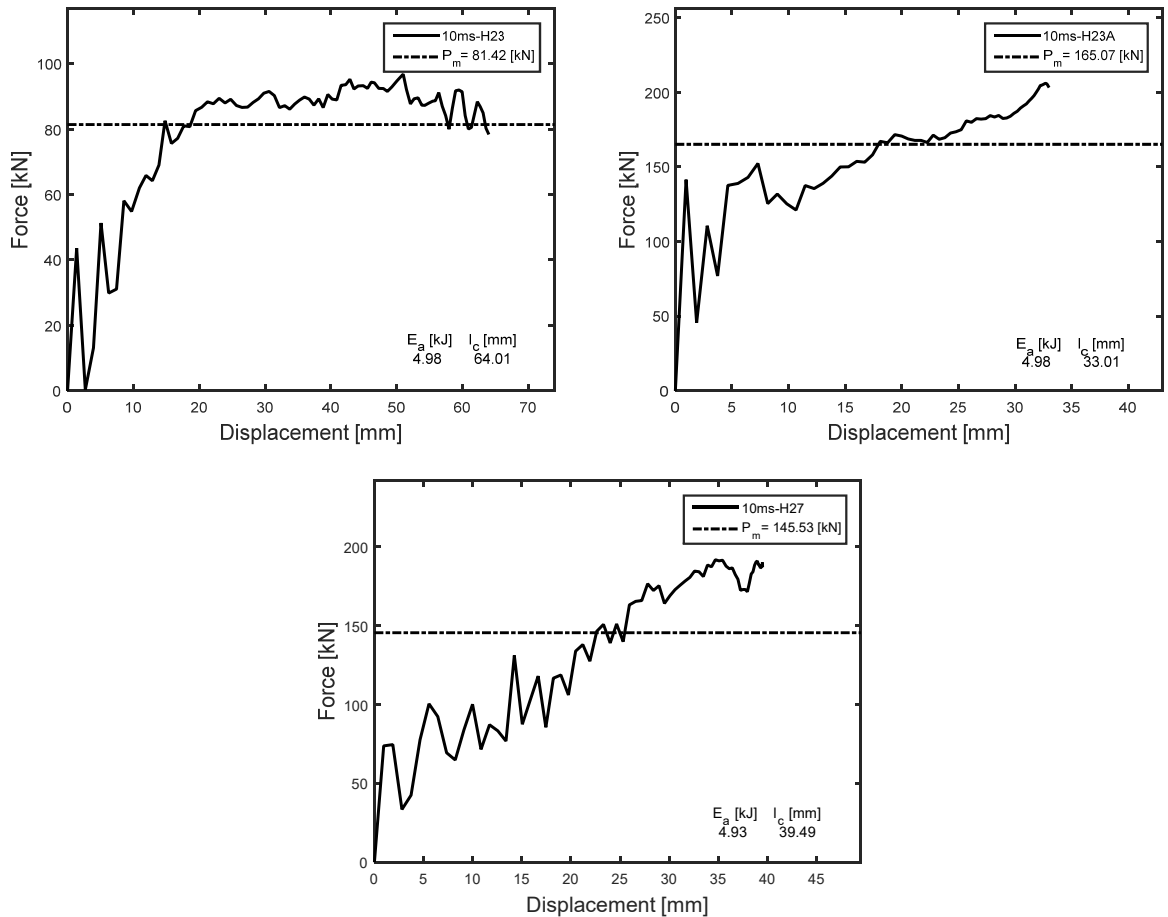


Figure 6: Load-displacement responses for $V=10\text{m/s}$.

Figure 6 shows the load-displacement responses for a speed $V=10\text{ m/s}$. It is observed that the impact speed affects the maximum crushing length of the three tubes. It is obvious that an increasing speed represents an increase in the kinetic energy to be absorbed. Concentric tubes with the smaller mesh H23 achieves the largest crushing length ($l_c= 64.01\text{ mm}$), the tube is highly deformed since there is not enough material for plastic deformation to absorb the kinetic energy of the mass. The shortest crushing length is attained for tube with mesh H23A ($l_c= 33.01\text{ mm}$) with larger cross-section than that of mesh H23.

Figure 7 and 8 show the load-displacement responses for speeds of $V=20\text{ m/s}$ and $V=30\text{ m/s}$, respectively. There is an increasing crushing length for an increasing speed. Once again the larger

crushing length is attained by the tube with the smallest mesh H23. Computed values for the crushing length are also presented in Table 3.

All responses exhibit fluctuations due to the dynamic behavior of the material. However when the displacements attain the crushing length l_c , the tube is almost compacted and the compression space is utilized, giving a rise in the crushing force.

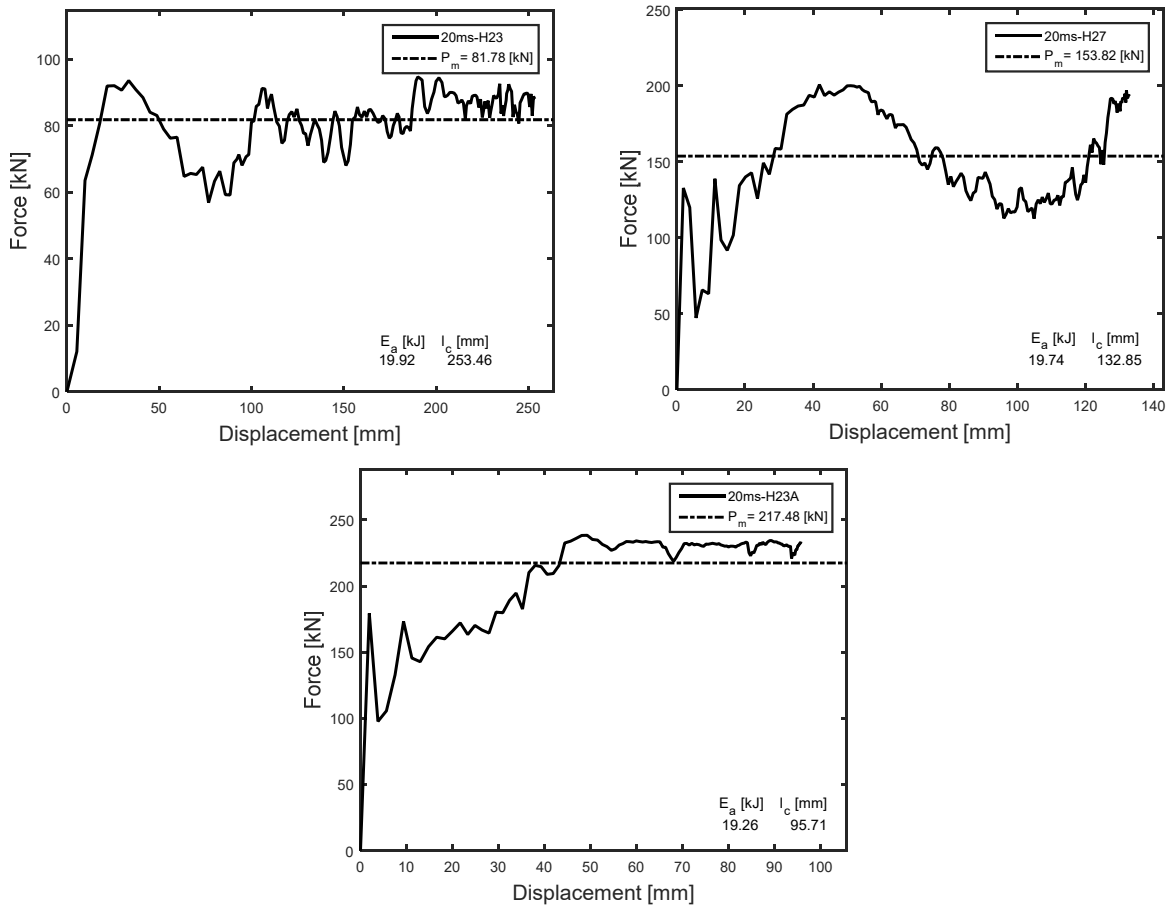


Figure 7: Load-displacement responses for $V=20\text{m/s}$.

Tube	V (m/s)	l_c (mm)	P_m (kN)	E_a (kJ)	l_c / L (%)
H23	10	64.01	81.42	4.98	15.65
	20	253.46	81.78	19.92	61.97
	30	346.74	97.13	31.54	84.78
H27A	10	39.49	145.53	4.93	9.70
	20	132.85	153.82	19.74	32.64
	30	259.58	189.42	44.82	63.78
H23A	10	33.01	165.07	4.98	7.71
	20	95.71	217.48	19.26	22.36
	30	194.17	250.16	45.01	45.37

Table 3: Numerical results for the concentric tubes.

As seen in Table 3, all tubes are capable to fully absorb the kinetic energy for $V=10$ m/s ($E_k= 5$ kJ) and $V=20$ m/s ($E_k= 20$ kJ). For $V= 30$ m/s, tube H23A is the only one capable to absorb the impact ($E_k= 45$ kJ) with a reasonable crushing length. For tube H23, it is physically impossible to withstand the impact at $V=30$ m/s, the absorbed energy was $E_a= 31.54$ kJ since the required crushing length is larger than its actual length.

The differences in the crushing length is mainly attributed to the amount of material in the cross section area. H27 and H23A tubes have more strand thickness than H23 tubes, but H27 has better cell aspect ratio than H23A.

On the other hand, the impact speed has a small effect in the mean load due to the hardening effect of the strain velocity as pointed in Eq. (3), but the crushing length is increased to achieve the total energy.

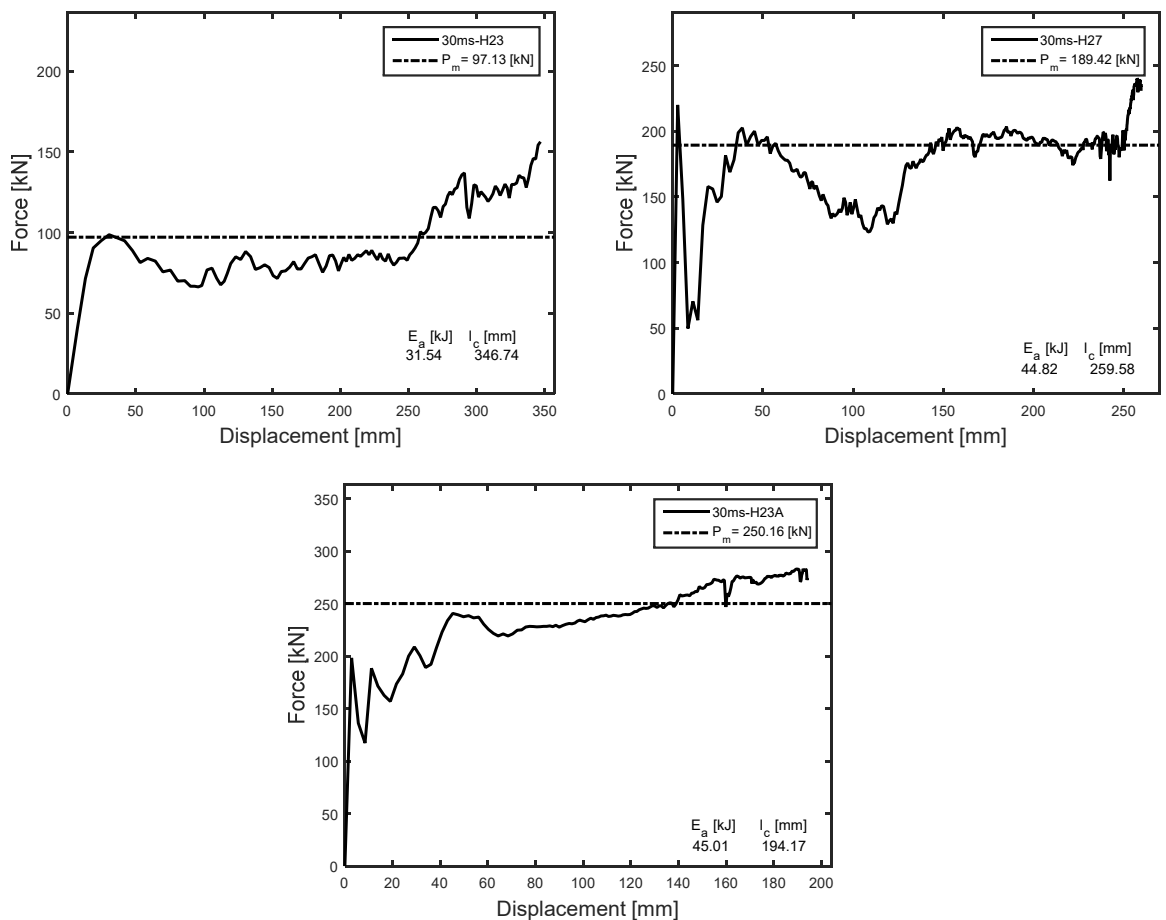


Figure 8: Load-displacement responses for $V=30$ m/s.

Analyzing the collapse of the concentric tubes (Figures 9 and 10), it was observed that the wall of the inner tube ($\alpha = 90^\circ$) moved outwards, until getting in contact with the wall of the outer tube

($\alpha = 0^\circ$). As both tubes collapsed, the wall of the outer tube limits the radial displacement of the wall in the inner, restricting its mode of collapse inducing additional plastic deformation in the inner.

Additionally, it is worth noticing that the concentric tube H23A exhibited a more controlled collapse mode than tubes H23 and H27A. As indicated in Table 1, the length of the axes in mesh H23A are smaller than for the other two tubes, hence the outer tube is more effective controlling the outward displacement of the inner. On the contrary, for tubes H23 and H27A the cells are larger allowing more interaction between the walls of the tubes, consequently the deformation is more irregular.

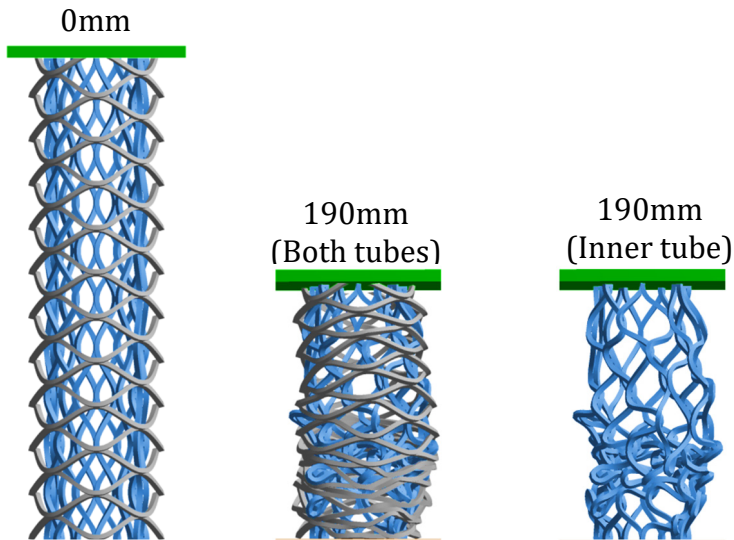


Figure 9: Deformed shape for Tube H23 ($V=30\text{m/s}$).

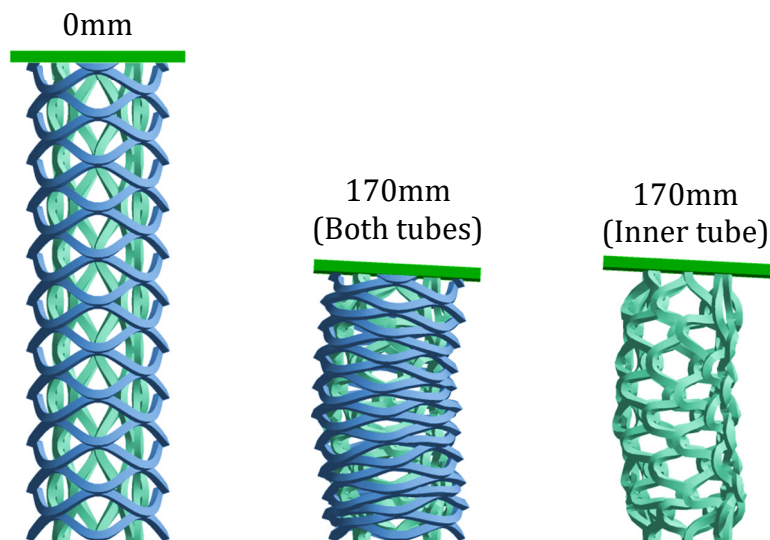


Figure 10: Deformed shape for Tube H23A ($V=30\text{m/s}$).

4 CONCLUSIONS

A numerical study on the axial crushing of expanded metal tubes subjected to impact load was conducted herein. During the axial collapse of concentric tubes 90 - 0, the inner tube with $\alpha = 90^\circ$ begins to fail by global buckling moving towards the wall of the outer tube with $\alpha = 0^\circ$, therefore the inner tube is restrained and forced to deform to fit a reduced volume. In this manner, the amount of energy absorbed by plastic deformation increases and the failure mode to collapse the tube is controlled, thus maintaining a constant collapse strength, which is a desirable property in energy absorption devices. Furthermore, a better cell aspect ratio enhances the energy absorption capacity for similar thickness, as can be seen in the mean force results. Finally, the impact speed has a small effect in the mean load due to the hardening effect of the strain velocity, but the crushing length is increased to absorb the total kinetic energy.

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